



## Article

# Stabilization of $S_3O_4$ at high pressure: implications for the sulfur-excess paradox

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## ABSTRACT

The amount of sulfur in  $SO_2$  discharged in volcanic eruptions exceeds that available for degassing from the erupted magma. This geological conundrum, known as the “sulfur excess”, has been the subject of considerable interests but remains an open question. Here, in a systematic computational investigation of sulfur-oxygen compounds under pressure, a hitherto unknown  $S_3O_4$  compound containing a mixture of sulfur oxidation states +II and +IV is predicted to be stable at pressures above 79 GPa. We speculate that  $S_3O_4$  may be produced via redox reactions involving subducted S-bearing minerals (e.g., sulfates and sulfides) with iron and goethite under high-pressure conditions of the deep lower mantle, decomposing to  $SO_2$  and S at shallow depths.  $S_3O_4$  may thus be a key intermediate in promoting decomposition of sulfates to release  $SO_2$ , offering an alternative source of excess sulfur released during explosive eruptions. These findings provide a possible resolution of the “excess sulfur degassing” paradox and a viable mechanism for the exchange of S between Earth’s surface and the lower mantle in the deep sulfur cycle.

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## 1. Introduction

Sulfur (S) is one of the major multi-valent volatile elements distributed broadly throughout the Earth, participating in a variety of fundamental geochemical processes (e.g., global biochemical circulation [1], metal transport [2], atmospheric S loading during volcanic eruptions, and core–mantle segregation [3], etc.). The chemical speciation of S is strongly influenced by the wide range of oxidation states available. Under highly reducing environments, it dominantly exhibits an oxidation state of –II as sulfide, whereas under strongly oxidizing conditions, it has an oxidation state of +VI in sulfate. Other chemical species where S takes up intermediate oxidation states, such as polysulfides, elemental S, sulfite, or thio-sulfate sulfite, may also exist in different geochemical settings [2,4,5]. The behavior of S in natural processes associated with complex oxidation–reduction reactions is unpredictable due to these

variable oxidation states across the range of –II to +VI. Therefore, the geochemical behavior of S is thus replete with paradoxes, and there are many open questions concerning geochemical processes related to S-bearing minerals.

A well-known geological paradox, “sulfur excess degassing”, has been recorded at numerous subduction zone volcanoes [6,7], where the amount of S (principally in the form of  $SO_2$ ) released during explosive eruptions may be orders of magnitude higher than that estimated for degassing of the erupted melt [5]. Various sources for the excess S released by magmas in volcanic emissions [8,9] have been proposed, including dissolution in silicate liquid [10,11] or a coexisting gas phase at depth before eruption [7,12], gas expulsion from magma mixing [13,14], crystallization-induced exsolution (second boiling) [15], or the breakdown of S-bearing minerals [16], etc. These mechanisms are based on magmatic systems related to volcanic eruptions in shallow crust. However, the ultimate source of S found near the Earth’s surface is derived from the Earth’s mantle [17]. Oxygen, as one of the most abundant elements, provides critical control on the nature of Earth S reservoirs. Compounds formed by S and O have important implications for geochemical processes and the nature of these

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S reservoirs. Key questions that need to be resolved regards the formation and properties of S–O compounds under mantle conditions.

Various S–O compounds such as  $\text{SO}_2$  [18],  $\text{SO}_3$  [19,20],  $\text{S}_7\text{O}$  [21], and  $\text{S}_8\text{O}$  [22] have been proposed experimentally at ambient pressure. However, the high pressures characteristic of the mantle may strongly modify the chemical properties of elements and promote the formation of unexpected minerals [23–26]. Several high-pressure  $\text{SO}_3$  phases have been proposed theoretically [27], but only  $\text{SO}_2$  has been studied experimentally at pressures of up to 60 GPa [28], and S–O compounds are not well understood at high pressures. A pressing task is therefore to investigate the viability of S–O compounds under pressure conditions relevant to the Earth's mantle.

Here, we report an extensive exploration of high-pressure phase diagrams of S–O compounds. Besides known  $\text{SO}_2$  and  $\text{SO}_3$  compounds, an unexpected stoichiometry of  $\text{S}_3\text{O}_4$  with an intriguing crystal structure, containing a mixture of +II and +IV S oxidation states, is predicted to appear at high pressures. We show that  $\text{S}_3\text{O}_4$  is produced in reactions of sulfates and sulfides with iron and goethite under high pressure conditions in the deep mantle, decomposing to  $\text{SO}_2$  and S under low  $P$ - $T$  conditions at shallow depths of the Earth, thus offering insight into S cycles, and the origin of excess S degassing in volcanic eruptions.

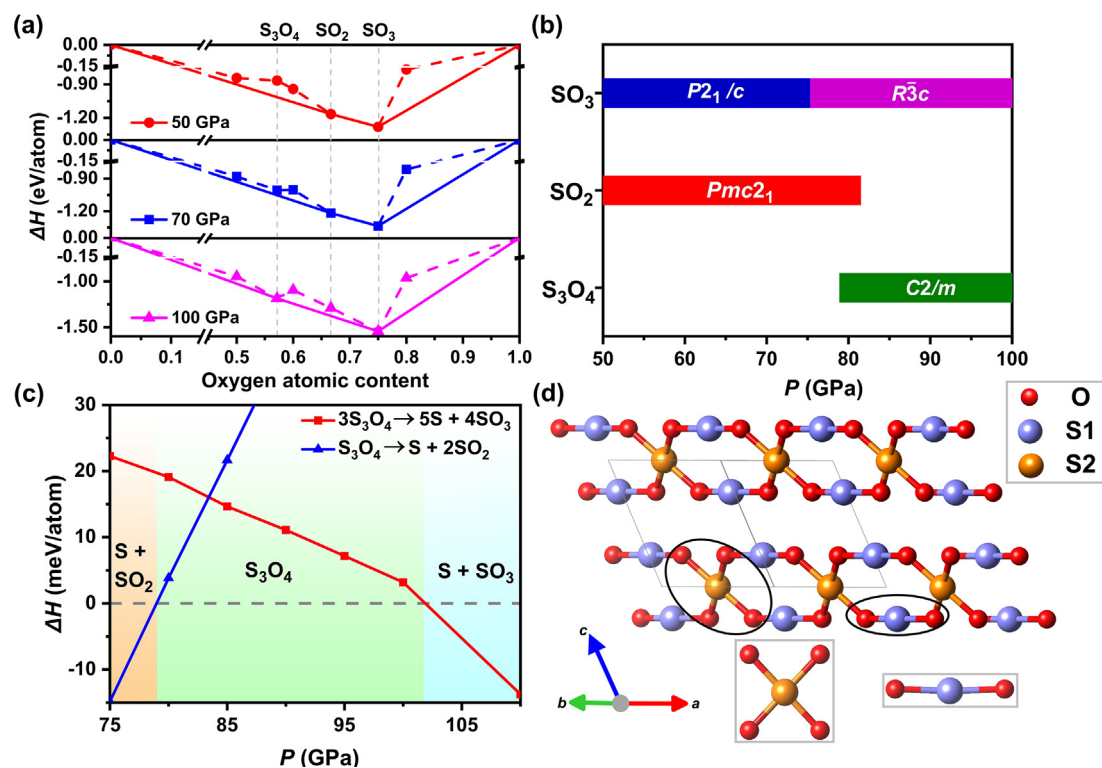
## 2. Methods

Crystal structure searches on  $\text{S}_x\text{O}_y$  ( $x = 1-3$ ,  $y = 1-4$ ) at pressures of 50, 70, and 100 GPa were undertaken using the swarm intelligence based-CALYPSO method [29–31], which has successfully resolved crystal structures of many materials at high pressures [32]. Note that the maximum simulation cell for structure searches contains 40 atoms for each composition. Structural optimization,

electronic structure, and phonon calculations were performed in a framework of density functional theory within the generalized gradient approximation [33] as implemented in the Vienna *ab initio* simulation package (VASP) [34]. Electron-ion interaction was described by projector augmented-wave potentials [35], with  $3s^23p^4$  and  $2s^22p^4$  configurations treated as the valence electrons of S and O, respectively. We also performed the full-potential all-electron calculations for the equation of state for  $\text{SO}_3$  over the considered pressure range using the WIEN2k code [36]. VASP results were near identical to those of the all-electron calculations (Fig. S1 online), validating the accuracy of the pseudopotentials. A kinetic cut-off energy of 900 eV and a spacing of  $2\pi \times 0.03 \text{ \AA}^{-1}$  for Monkhorst-Pack  $k$ -mesh sampling [37] were adopted to give converged total energies ( $\sim 1$  meV/atom). Ionic positions were fully relaxed until the residual force acting on each ion was less than 1 meV/Å. Due to the layered structure of  $\text{S}_3\text{O}_4$ , the influence of van der Waals (vdW) interactions was considered using the optB88-vdW functional [38]. To describe the localized 3d electrons of Fe atoms in reactants, we considered electron correlation by using the GGA + U [39] method, with on-site Coulomb interaction of  $U = 5.0$  eV and a Hund coupling constant of  $J = 0.8$  eV [40]. The dynamic stabilities of the predicted new phase of  $\text{S}_3\text{O}_4$  were verified by phonon calculations using the Hellmann-Feynman theorem with the finite displacement method of the  $2 \times 2 \times 2$  supercells containing 112 atoms, as implemented in the PHONOPY code [41,42].

## 3. Results

The main structure search results are depicted in convex hull diagrams in Fig. 1a. The energetic stabilities of various S–O structures were evaluated from their formation enthalpies relative to



**Fig. 1.** Relative thermodynamic stability of the S–O system at 0 K. (a) Convex hull data of the  $\text{S}_{1-x}\text{O}_x$  system at 50, 70, and 100 GPa. Formation enthalpies,  $\Delta H$ , for each structure were calculated with respect to elemental S and O solids by  $\Delta H(\text{S}_{1-x}\text{O}_x) = H(\text{S}_{1-x}\text{O}_x) - (1-x)H(\text{S solid}) - xH(\text{O solid})$  ( $0 < x < 1$ ). Known S-III [43], S-IV [44], and  $\epsilon\text{-O}_2$  [45] phases were selected as the reference structures in the corresponding stable pressure ranges. Stable structures are located on the solid lines, and metastable structures sit on the dashed lines. (b) Predicted pressure-composition phase diagram of S–O phases. (c) Calculated pressure-enthalpy diagram for the reactions  $3\text{S}_3\text{O}_4 \rightarrow 4\text{SO}_3 + 5\text{S}$  and  $\text{S}_3\text{O}_4 \rightarrow 2\text{SO}_2 + \text{S}$  using optB88-vdW functional. The zero-point energy was included in the above energy calculations. (d) Crystal structure of  $C2/m$ - $\text{S}_3\text{O}_4$  containing mixed two-fold (S1) and four-fold (S2) coordination of S.

dissociation products of the relevant elemental S [43,44] and O solids [45]. At pressures of 50 and 70 GPa, the known stoichiometries  $\text{SO}_2$  and  $\text{SO}_3$  were readily identified as being stable in the structure search simulations.  $\text{SO}_3$  was found to be the most stable phase against decomposition throughout the studied pressure range (50–100 GPa). At 100 GPa, an unexpected composition of  $\text{S}_3\text{O}_4$  became stable with respect to the dissociation products of elemental S and  $\text{SO}_3$  (Fig. S2 online). More detailed structural information can be found in Table S1 (online). Predicted stable pressure ranges for the considered structures are listed in Fig. 1b.  $\text{S}_3\text{O}_4$  is energetically favorable relative to decomposition into elemental S and  $\text{SO}_2$  or  $\text{SO}_3$  in the pressure range of 79–102 GPa (Fig. 1c). The emergence of  $\text{S}_3\text{O}_4$  leads to the instability of  $\text{SO}_2$  above 81.5 GPa. The different exchange–correlation functionals and vdW methods have also demonstrated the stability of  $\text{S}_3\text{O}_4$  at high pressure and temperature ( $P$ - $T$ ) (Figs. S3 and S4 online). We calculated phonon dispersions and observed no imaginary frequencies for the  $\text{S}_3\text{O}_4$  structures at 100 GPa (Fig. S5 online), indicating that the predicted structure is dynamically stable.

The structure of  $\text{S}_3\text{O}_4$  (Fig. 1d) is inherently layered and contains mixed two- and four-fold S coordination. Specifically, S1 is linearly coordinated to two O atoms, and S2 is square-coordinated to four O atoms. All S atoms are bonded to two adjacent S atoms, thus forming zigzag polymeric all-S chains. The S1–S1 and S1–S2 bond lengths are 2.22 and 2.13 Å at 80 GPa, respectively, slightly longer than the S–S bond lengths (2.01 Å) in the S-III phase, therefore indicating relatively weaker covalent S–S bonding. To further elucidate the nature of the bonding, we have examined the electron localization function (ELF) [46] of  $\text{S}_3\text{O}_4$  in the (100) and (010) planes (Fig. 2a). Two inequivalent S atoms were clearly seen, with a less localized charge distribution is seen on the S–O bonds, indicating a significant degree of ionicity between the O anions and S cations. Clear covalent S–S bonding is indicated by the strong charge localization between the nearest neighbor S–S.

The oxidation states of S in geological environments play pivotal roles in planetary chemical and physical dynamics [49]. In general, the oxidation state of an element is closely related to local coordination and charge transfer. The S oxidation states in  $\text{SO}_2$  and  $\text{SO}_3$  can be assigned unambiguously as +IV and +VI, respectively. In contrast, the two- and four-fold coordination of S atoms with O atoms in  $\text{S}_3\text{O}_4$  reveals its mixed-valence state. A Bader charge analysis [50], summarized in Table 1, supports this interpretation. The Bader charges systematically underestimate the formal charge state ( $\text{O}^{2-}$  here has a charge  $-1.28e$  in the  $\text{SO}_2$ ). In  $\text{SO}_3$ , S has a formal charge state of +VI, with a charge transfer of  $3.90e$  from S to O, similar to that in  $\text{SF}_6$  ( $\sim 3.73e$ ). In  $\text{SO}_2$ , S has a formal charge state of +IV, with a charge transfer is  $2.56e$ . In  $\text{S}_3\text{O}_4$ , the partial charge of  $2.68e$  in square-coordinated S2 almost equals that of the  $\text{SO}_2$  case, so S2 can be considered as having an oxidation state of +IV. However, S1 is significantly less positively charged ( $1.04e$ ) than the  $\text{S}^{+4}$

**Table 1**

Partial charges for various S–O compounds obtained from Bader integration at 80 GPa.

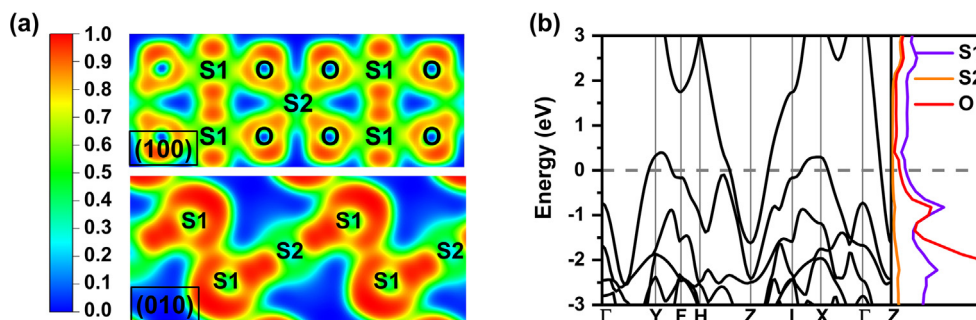
Compounds	S (e)		O (e)
	S1	S2	
$\text{S}_3\text{O}_4$	+1.04	+2.68	−1.19
$\text{SO}_2$		+2.56	−1.28
$\text{SO}_3$		+3.90	−1.30

anion in  $\text{SO}_2$ . This highlights a crucial distinction of the S1 from S in  $\text{SO}_2$ , indicating that the linearly coordinated S1 in  $\text{S}_3\text{O}_4$  adopts the rare +II S oxidation state.

S–O compounds tend to be insulating, with satisfaction of the octet rule usually leading to the opening of a band gap. This rule is applicable to the predicted polymeric phases of  $\text{SO}_2$  and  $\text{SO}_3$ . However, in  $\text{S}_3\text{O}_4$ , two bands were found to cross the Fermi level, forming an electron pocket around the Z point and a hole pocket spanning the X and Y points (Fig. 2b), giving rise to a clear metallic character of  $\text{S}_3\text{O}_4$ . The projected density of states (Fig. 2b) indicates that both O and linearly coordinated S1 contribute to the density of electronic states at the Fermi level, with the latter contribution being dominant. The metallic character originates from an overlap of the S1 electron lone pairs, which depends on the interlayer distance (Fig. S6 online).

Both S and O are typical light elements, so the stability of S–O compounds may be sensitive to temperature. To assess viability at high temperature, we further examined their energetic and structural stability under high  $P$ - $T$  conditions. The free energies, including vibrational contributions and entropic effects, were evaluated for each phase using the quasi-harmonic approximation [51]. Formation enthalpy calculations further indicate that  $\text{S}_3\text{O}_4$  is energetically favorable relative to decomposition into  $\text{SO}_3$  and S above 70 GPa, with temperature having a minor effect on threshold pressure (Fig. 3a). Against decomposition into  $\text{SO}_2$  and S, the stability region of  $\text{S}_3\text{O}_4$  shifted to higher pressures with increasing temperature, from 79 GPa at 0 K to 100 GPa at 2300 K (Fig. 3b). *Ab-initio* molecular dynamics calculations indicate that  $\text{S}_3\text{O}_4$  remains firmly solid at 2000 K at pressures of 80–100 GPa, corresponding to deep mantle conditions (Fig. S7 online), indicating that  $\text{S}_3\text{O}_4$  may exist in solid form in the deep mantle. Overall, the predicted  $\text{S}_3\text{O}_4$  is stable under  $P$ - $T$  conditions relevant to the Earth's lower mantle [52], but decomposes into  $\text{SO}_2$  and S at low pressure.

It is well known that the exchange of S between the Earth's surface and mantle, involving transport of S to the mantle via subduction and return to the surface by volcanic degassing, results in a global S cycle [53]. Nearly 43.6 Tg per year of S in the form of such as sulfates (e.g.,  $\text{CaSO}_4$  and  $\text{MgSO}_4$  [53]) and sulfides (e.g., FeS and  $\text{FeS}_2$  [54,55]) are subducted into the deep mantle [17]. However, the average annual  $\text{SO}_2$  flux from erupting volcanoes is estimated to be 11.9 Tg (1 Tg =  $10^{12}$  g) [56]. The cycle involves transformations



**Fig. 2.** The electronic properties of  $\text{S}_3\text{O}_4$ . (a) Calculated ELF in the (100) and (010) planes of  $\text{C2}/m\text{-S}_3\text{O}_4$  at 80 GPa. (b) Band structures and projected density of states of  $\text{C2}/m\text{-S}_3\text{O}_4$  at 80 GPa. Dashed line indicates the Fermi level. Band structures were calculated using the Heyd–Scuseria–Ernzerhof hybrid functional [47,48].

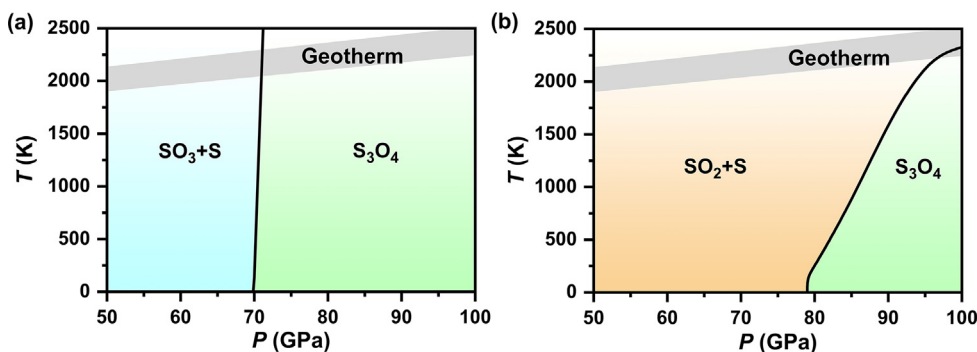


Fig. 3. The *P-T* phase diagrams for the reactions of  $4\text{SO}_3 + 5\text{S} \rightarrow 3\text{S}_3\text{O}_4$  (a) and  $2\text{SO}_2 + \text{S} \rightarrow \text{S}_3\text{O}_4$  (b). The geotherm curve was adapted from Ref. [52]

of S species via redox-driven chemical processes such as sulfate reduction and sulfide oxidation. It is estimated that ~1 wt% metallic Fe is present due to self-reduction reactions in the lower mantle [57]. In addition, the pyrite-type FeOOH may survive under the *P-T* conditions of the lower mantle within deeply subducted slabs

[58]. Therefore, we explored the possibility of  $\text{S}_3\text{O}_4$  production through decomposition of sulfates with Fe as a reducing agent or through oxidation reactions of sulfides with FeOOH as an oxidizing agent. We explored 39 possible reaction routes (Fig. 4a), in which iron oxides and hydrogen-bearing minerals that may exist in the

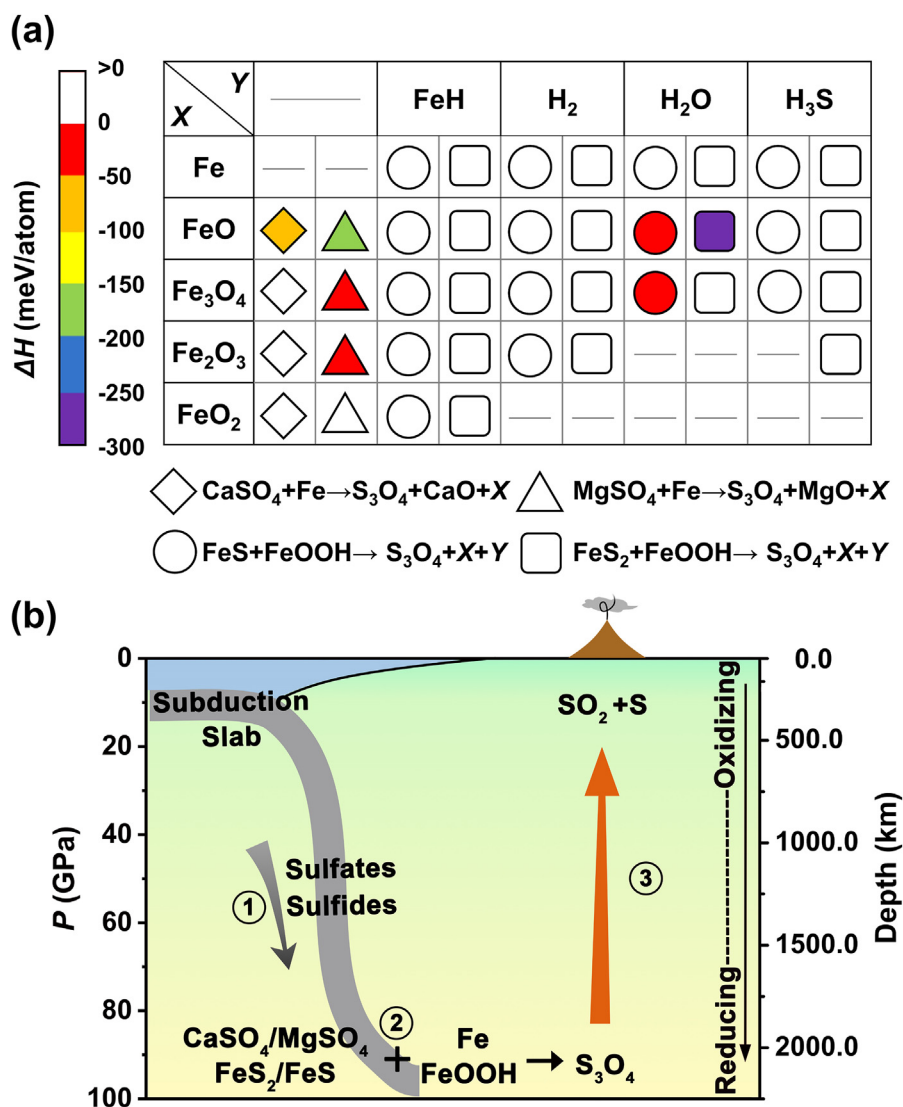


Fig. 4. Phase equilibria and the proposed S cycle in Earth. (a) Relative enthalpy of proposed reactions forming  $\text{S}_3\text{O}_4$  at 100 GPa. Rhombus and triangle markers represent the reactions of  $\text{CaSO}_4$  and  $\text{MgSO}_4$  with iron; circles and squares represent the reactions of FeS and  $\text{FeS}_2$  with FeOOH. X and Y represent different iron-oxides and hydrogen-bearing compounds, respectively. Colors denote enthalpies for the corresponding reactions. Crystal structures adopted to evaluate enthalpies are presented in Table S2 (online). (b) Processes proposed for the exchange of S between the Earth's surface and mantle: ① sulfates and sulfides are subducted into the deep mantle; ② subducted sulfates and sulfides react with the Fe or FeOOH to produce  $\text{S}_3\text{O}_4$ ; and ③  $\text{S}_3\text{O}_4$  decomposes to  $\text{SO}_2$  and elemental S, with emission of  $\text{SO}_2$  to the surface via volcanism.

Earth's interior were chosen as products. The calculated negative reaction enthalpies of seven routes support the formation of  $S_3O_4$  via redox reactions at 100 GPa, reflecting deep-mantle pressure. The pressure dependence of enthalpies of these reactions is shown in Fig. S8 (online). According to our calculations,  $S_3O_4$  can be produced by two almost opposing mechanisms (i.e., reduction of sulfates and oxidation of sulfides), with these reactions being strongly dependent on redox conditions. It is well known that redox conditions have evolved continuously in the Earth's history [59]. Therefore,  $S_3O_4$  compounds may have been produced by redox reactions involving S-bearing minerals throughout geological history.

Our results indicate three possible processes for S cycling in the Earth (Fig. 4b). Firstly, S-bearing sulfates or sulfides (e.g.,  $CaSO_4$ ,  $FeS_2$  [60], and  $FeS$  [61]) are transported to the deep mantle in subduction slabs. They react with Fe or  $FeOOH$  (present in the mantle) to produce  $S_3O_4$  under reducing or oxidizing conditions. If  $S_3O_4$  formed in the deep mantle ascends (through mantle dynamic processes) to shallow depths of the Earth with low pressure conditions, it decomposes to S and  $SO_2$ , with the latter being the main form of S released during explosive eruptions (Fig. 3b). The mechanism of direct decomposition of S-bearing minerals (e.g.,  $CaSO_4$ ,  $FeS_2$ , and  $FeS$ ) to release  $SO_2$  as an explanation of the sulfur-excess paradox is not supported, even at high pressure [11] (see Fig. S9 online for enthalpy calculations); however,  $S_3O_4$ , which has not been considered previously, provides an alternative S reservoir in the deep mantle, completing the deep S cycle and aiding explanation of the paradox in volcanic eruptions, especially in some hotspots and for mantle plume volcanism (e.g., Nyamuragira volcano) [62].

#### 4. Conclusion

A hitherto unknown compound,  $S_3O_4$ , has been identified as being stable under the high *P-T* conditions of the deep mantle. It contains a mixture of S(II) and S(IV) oxidation states and exhibits a peculiar metallic nature. A systematic examination of formation and decomposition reactions indicates that  $S_3O_4$  may be a key factor in promoting redox reactions of sulfate or sulfide in the deep mantle and in the release of  $SO_2$  at shallow depths, thereby offering insight into the origin of excess S degassing observed in volcanic eruptions. These results are of fundamental significance and have implications for processes in chemistry and geoscience; further experimental exploration is expected.

#### Conflict of interest

The authors declare that they have no conflict of interest.

#### Acknowledgments

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#### Author contributions

Yanchao Wang conceived and designed the project; Siyu Liu, Pengyue Gao, and Yanchao Wang performed simulations; Siyu Liu, Pengyue Gao, Andreas Hermann, Guochun Yang, Jian Lü, Yanming Ma, Ho-Kwang Mao, and Yanchao Wang performed data analysis and interpretation of the results; Siyu Liu, Pengyue Gao,

Jian Lü, Yanming Ma, Ho-Kwang Mao, and Yanchao Wang wrote the paper; and all authors contributed to discussions of the results and revisions of the manuscript.

#### Appendix A. Supplementary materials

Supplementary materials to this article can be found online at <https://doi.org/10.1016/j.scib.2022.01.005>.

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