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# Engineering mixed polyanion red-emitting $Rb_2Bi(PO_4)(WO_4):Eu^{3+}$ phosphors with negligible thermal quenching and high quantum yield



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## ABSTRACT

The high-performance red-emitting phosphors are strongly demanded as the next generation solid state lighting materials. Benefitting from the favorable accommodation capacity of polyanions and high charge transfer energy state in the structure, the pure red-emitting Rb<sub>2</sub>Bi(PO<sub>4</sub>)(WO<sub>4</sub>):Eu<sup>3+</sup> (RBPWO:Eu<sup>3</sup> materials with negligible thermal quenching and high quantum yield were synthesized. RBPWO:Eu<sup>3+</sup> has high absolute quantum yields of 88.1% and exceptional thermal stability of 98.6% integral retention at 423 K relative to the initial value at 298 K. Additionally, affected by the low symmetry crystalline-field of  $D_{2h}$ , the title phosphors exhibit intense absorption band peaking at 394 nm and pure red light with the color CIE values (0.6482, 0.3515). The white light-emitting diode (w-LED) device fabricated from the title material exhibits low correlated color temperature (4000 K), high color rendering index ( $R_a = 90$ ) and moderate luminous efficiency (18.07 lm/W) in the preliminary experiment. The present work indicates that the RBPWO:Eu<sup>3+</sup> phosphors are promising red components for the near ultraviolet light excited w-**I FDs** 

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## 1. Introduction

White light-emitting diodes (w-LEDs) are widely used in indoor and outdoor illumination due to the low energy consumption, high brightness, long working life and environmental friendliness [1-3]. The popular implementation approach for fabricating commercial w-LEDs is coating the yellow  $Y_3Al_5O_{12}$ :Ce<sup>3+</sup> (YAG:Ce<sup>3+</sup>) materials on the blue-emitting InGaN chips [4-8]. However, they suffer from an insufficient color rendering index (R<sub>a</sub>) and high correlated color temperature  $(T_c)$  [9]. To reach the high-quality white light, a promising solution was developed: utilizing the near ultraviolet (NUV) light to pump the three primary colors (blue, green and red) phosphors [10–13]. The driver for such applications has involved the extensive exploration for high-performance red-emitting materials matched with the NUV light emitting chips.

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To date, several commercial red phosphors including  $Sr_2Si_5N_8:Eu^{2+}$  [14], (Sr,Ca)AlSiN\_3:Eu^{2+} [15],  $Y_2O_3:Eu^{3+}$  [16] and  $Y_2O_2S:Eu^{3+}$  [17] have been developed. Although the Sr<sub>2</sub>Si<sub>5</sub>N<sub>8</sub>:Eu<sup>2+</sup> [14] and (Sr,Ca)AlSiN<sub>3</sub>:Eu<sup>2+</sup> [15] materials possess broad band emission, their luminous efficacies are influenced due to a portion of the spectra located in the eye-insensitive region (above 700 nm). Additionally, the synthesis of bivalence europium based materials requires a reducing atmosphere, which results in the higher production costs [18-20]. It seems that the Eu<sup>3+</sup> ion is a promising activator because of its narrow red line emission. However, the Y<sub>2</sub>O<sub>2</sub>S:Eu<sup>3+</sup> materials are easily deliquescence to release toxic sulfide gases and their service lives are reduced. In general, metal oxide compounds have excellent thermal stability. For example, the melting point of  $Y_2O_3$  is over 2000 °C. The Eu<sup>3+</sup> doped  $Y_2O_3$ phosphors have intense absorbance centred at 255 nm in UV region, but they can't effectively be excited by NUV light [21]. Among the emission bands of Eu<sup>3+</sup>, only the  ${}^{5}D_{0}$ - ${}^{7}F_{2}$  transition is responsible for the typical red luminescence emission. In view of the luminescence behaviors of Eu<sup>3+</sup> are susceptible to their crystallographic environments, designing good structure of the matrix is

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vital [22–25]. First of all, the excellent accommodation capacity for Eu<sup>3+</sup> ions is essential. Based on the Vegard Law, if the radii values difference of two metal ions is within 15%, they can be replaced each other easily [26]. Among the eight-coordinated trivalent metal ions, including  $Sc^{3+}$ ,  $Y^{3+}$ ,  $La^{3+}$ ,  $Gd^{3+}$ ,  $Lu^{3+}$  and  $Bi^{3+}$ , the radius difference of  $Gd^{3+}$  or  $Bi^{3+}$  and  $Eu^{3+}$  is less than five percent. Therefore, the gadolinium or bismuth based compounds are considered as suitable host compounds [27]. In addition, the existence of concentration quenching effect results in the limited content of activators ions, which will lead to insufficient fluorescence quantum yields [28,29]. For instance, the critical quenching contents ( $x_c$ ) of Eu<sup>3+</sup> in CaGdAlO<sub>4</sub>:Eu<sup>3+</sup> and Ba<sub>2</sub>YAlO<sub>5</sub>:Eu<sup>3+</sup> are 7 mol % and 30 mol %, respectively [29]. Some large anionic groups such as PO<sub>4</sub> and WO<sub>4</sub>, can segregate the adjacent luminescence activators that contribute to reducing the interaction between adjacent europium ions and increasing luminous particles density. The  $x_c$  values in Ca<sub>9</sub>Bi(PO<sub>4</sub>)<sub>7</sub>:Eu<sup>3+</sup> and NaSrLa(WO<sub>4</sub>)<sub>3</sub>:Eu<sup>3+</sup> are 60 and 80 mol %, respectively [30,31]. Based on Binnemans theory, the <sup>5</sup>D<sub>0</sub>-<sup>7</sup>F<sub>2</sub> transition of Eu<sup>3+</sup> is hypersensitive for the crystallographic surroundings and the low symmetry host structure is beneficial for strengthening the <sup>5</sup>D<sub>0</sub>-<sup>7</sup>F<sub>2</sub> transition [24,32].

Herein, the  $Rb_2Bi(PO_4)(WO_4):Eu^{3+}$  (RBPWO:Eu<sup>3+</sup>) materials were investigated. In the structure of RBPWO: $Eu^{3+}$ , the mixed polyanions  $PO_4$  and  $WO_4$  coexist and the  $Eu^{3+}$  ions can easily substitute  $Bi^{3+}$  at any  $Eu^{3+}/Bi^{3+}$  ratio, and the  $x_c$  is capable to reach 90%. Meanwhile, the high absolute quantum yields (AQY) of 88.1% was obtained. In addition, RBPWO belongs to orthorhombic system, thus its low symmetry structure of  $D_{2h}$  makes the  ${}^{5}D_{0}$ - ${}^{7}F_{2}$  transition of  $Eu^{3+}$  very intense and dominating the emission spectrum [32]. The thermal stability of trivalent europium doped materials relies heavily on the charge-transfer (CT) state of Eu<sup>3+</sup>. Blasse considered that the electrons can absorb thermal energy and then jump to the CT state. Afterwards, they return to ground state directly without emitting light [33]. Obviously, the phosphors with the high-level CT states should exhibit high quenching temperature and good thermal stability. In RBPWO:Eu<sup>3+</sup>, the highly positioned CT energy level (260 nm) of Eu<sup>3+</sup> was observed, and the exceptional thermal stability of the phosphors was obtained that is 98.6% integral retention at 423 K relative to the initial value at room temperature. Thereby, the red-emitting phosphors RBPWO:Eu<sup>3+</sup> have pure color, good thermal stability and high quantum efficiency. This high-efficiency red light emission phenomenon can be observed in other Rb<sub>2</sub>Bi(-PO<sub>4</sub>)(WO<sub>4</sub>)-type phosphors containing mixed polyanion, such as Na<sub>2</sub>Gd(PO<sub>4</sub>)(MoO<sub>4</sub>) and Na<sub>2</sub>Lu(PO<sub>4</sub>)(MoO<sub>4</sub>) [34-38]. Herein the site occupancy, optical properties, heat resistance properties, AQY and the potential application of RBPWO:Eu<sup>3+</sup> were systematically studied.

## 2. Experimental

## 2.1. Material synthesis

The Rb<sub>2</sub>Bi(PO<sub>4</sub>)(WO<sub>4</sub>):*x*Eu<sup>3+</sup> (RBPWO:*x*Eu<sup>3+</sup>) ( $0 \le x \le 1$ ) materials were prepared by the solid state reaction in the air atmosphere. The stoichiometric amounts of Rb<sub>2</sub>CO<sub>3</sub> (99%, Macklin Co., Ltd), Bi<sub>2</sub>O<sub>3</sub> (99.9%, Macklin Co., Ltd), NH<sub>4</sub>H<sub>2</sub>PO<sub>4</sub> (99%, Macklin Co., Ltd), WO<sub>3</sub> (99.5%, Macklin Co., Ltd) and Eu<sub>2</sub>O<sub>3</sub> (99.9%, Shanghai Yuelong Rare earth New Materials Co., Ltd) were mixed thoroughly and preheated at 500 °C for 12–36 h. Subsequently, the mixture was sintered at 700 °C for 48 h with 3–5 intermediate grindings.

#### 2.2. Characterization

The crystal structures were characterized with use of a Bruker D8 Advance X-ray diffractometer. The scanning step was 0.05 °s<sup>-1</sup>.

The structure was refined by the General Structure Analysis System (GSAS) and the EXPGUI software [39,40]. The morphology of the samples were observed by a field emission scanning electron microscope (MERLIN Compact, SEM). The optical diffuse reflectance spectra (DRS) were recorded on a Varian Cary 5000 spectrophotometer. The room-temperature and variable-temperature photoluminescence (PL), photoluminescence excitation (PLE) spectra and luminescence lifetimes were taken by an FLS-980 spectrofluorometer. The optical properties of fabricated *w*-LED device were measured by an Everfine LED instrument. The absolute quantum yields (AQY) were measured using an integrated sphere device according to the equation (1):

$$AQY = \frac{L_{Sample}}{E_{Reference} - E_{Sample}}$$
(1)

Where,  $L_{Sample}$  is the emission integrals of samples,  $E_{Reference}$  and  $E_{Sample}$  stand for the scatter integrals of the samples and the reference (a white tablet pressed by the BaSO<sub>4</sub> powders).

## 3. Results and discussion

#### 3.1. Rietveld refinement and phase identification

To check the structure of the as-synthesized RBPWO:xEu<sup>3+</sup> compounds, the crystal data RBPWO (ICSD# 426594) was adopted as an initial refinement model [32]. The observed and calculated XRD patterns for RBPWO and Eu<sup>3+</sup> completely substituted compound Rb<sub>2</sub>Eu(PO<sub>4</sub>)(WO<sub>4</sub>) are shown in Fig. 1. The final agreement factors converged to R<sub>p</sub> = 7.94% and R<sub>wp</sub> = 10.34% for RBPWO, R<sub>p</sub> = 7.86% and R<sub>wp</sub> = 9.94% for Rb<sub>2</sub>BiEu(PO<sub>4</sub>)(WO<sub>4</sub>), which verify the two compounds were synthesized successfully [41,42]. Additionally, in Fig. 2a, the diffraction curves are given for selected compositions and they match well with the pattern of matrix compound that verifies the successful incorporation of Eu<sup>3+</sup> ions into the host structure.

## 3.2. Micro-morphology

The morphologies of the RBPWO, RBPWO: $0.9Eu^{3+}$  and  $Rb_2Eu(-PO_4)(WO_4)$  compounds are presented in Fig. 3 and they are analogous. The samples comprised the plate-like particles with the size range of  $0.5-4 \mu m$ . The faceted plate-like particle morphology is an indicator of layered crystal structure, similar relation was observed in other crystal family of layered molybdates [43–45]. The strongly heterogeneity is stemming from the manual grinding after calcination. Meanwhile, the surface of the particles is clean and smooth, indicating the materials are highly crystallized and that is beneficial to realize the high light output.

#### 3.3. Crystal structure and preferable site occupancy

The RBPWO compound crystallizes in a centrosymmetric lattice with the space group of *l*bca. In the RBPWO structure, each cation (Rb<sup>+</sup>, Bi<sup>3+</sup>, P<sup>+5</sup> and W<sup>+6</sup>) has a unique crystallographic position [32]. Both Rb and Bi atoms are eight-coordinated to form the RbO<sub>8</sub> and BiO<sub>8</sub> dodecahedra, while P and W atoms are connected to four O atoms in forming the PO<sub>4</sub> and WO<sub>4</sub> tetrahedra (Fig. 4a). Along the *b* direction, BiO<sub>8</sub> and PO<sub>4</sub> polyhedron are combined to construct a two-dimensional  $\frac{2}{\infty}(Bi_2P_2O_{12})^{8-}$  chain (Fig. 4b). Further, the WO<sub>4</sub> tetrahedra are attached in both sides of the chain by sharing the O atoms of BiO<sub>8</sub> polyhedron, with Rb atoms occupying the voids between the neighbouring sheets.

As based on Shannon empirical relations, the ionic radii of eight-coordinated  $\rm Rb^+$ ,  $\rm Bi^{3+}$  and  $\rm Eu^{3+}$  are 1.60, 1.11 and 1.07 Å, and those



Fig. 1. Rietveld structure refinement plots for (a) RBPWO and (b) Rb<sub>2</sub>Eu(PO<sub>4</sub>)(WO<sub>4</sub>).



Fig. 2. (a) The representative XRD patterns, and (b) the magnified XRD patterns with the selected diffraction peaks from 19° to 20° of RBPWO:xEu<sup>3+</sup> ( $0 \le x \le 1$ ).

of four-coordinated  $P^{5+}$  and  $W^{6+}$  cations are 0.31 and 0.42 Å, respectively [27,46,47]. Apparently, compared to Eu<sup>3+</sup>, Rb<sup>+</sup> is too large, while  $P^{5+}$  and  $W^{6+}$  are too small, and only the size of  $Bi^{3+}$  is

very close to that of Eu<sup>3+</sup>. Therefore, the trivalent europium ions are mostly likely occupying the lattice sites of Bi<sup>3+</sup>. Besides, the large valence difference between the cations is further factor to exclude



Fig. 3. SEM images for (a) RBPWO, (b) RBPWO:0.9Eu<sup>3+</sup>, and (c) Rb<sub>2</sub>Eu(PO<sub>4</sub>)(WO<sub>4</sub>).



Fig. 4. (a) Coordination environments of cations and (b) crystal structure of RBPWO viewed along *a* axis.

the trivalent europium ion substitution for monovalent rubidium, pentavalent phosphorus or hexavalent tungsten ions. The representative diffraction peak of RBPWO: $xEu^{3+}$  ( $0 \le x \le 1.00$ ) appeared in the diffraction angle range from 19° to 20° is shown in Fig. 2b. Clearly, the characteristic peak (121) shifts gradually to higher angles with the Eu<sup>3+</sup> content increase. This is derived from the eight-coordinated Eu<sup>3+</sup> is slightly smaller than the same-coordinated Bi<sup>3+</sup> [27,46,47]. The decay curves of the RBPWO:Eu<sup>3+</sup> substances, as given in Fig. S1, match well with a single-exponential function and it further shows that Eu<sup>3+</sup> occupy only unique Bi<sup>3+</sup> crystallo-graphic site.

Based on Vegard Law, if the radii values difference of  $Bi^{3+}$  and  $Eu^{3+}$  cations is within 15%,  $Eu^{3+}$  can be dissolved into the host to form a solid solution [26]. Since the size difference of  $Bi^{3+}$  and  $Eu^{3+}$  cations is about 3.6%, the formation of wide-range RBPWO: $xEu^{3+}$  solid solutions can be predicted. As seen in Figs. 1 and 2, the  $Eu^{3+}$  ions can substitute  $Bi^{3+}$  ions at any  $Eu^{3+}/Bi^{3+}$  ratio without phase transition that is in good relation with the prediction. The structure of RBPWO is quite flexible for the incorporation of doped  $Eu^{3+}$  ions, and the RBPWO: $Eu^{3+}$  materials can be easily obtained in pure polycrystalline form.

The cell parameters of RBPWO: $xEu^{3+}$  compounds reduced linearly as the content of  $Eu^{3+}$  increased, as shown in Fig. 5 and Table S1. Intriguingly, the reduction rates of the cell parameters are different: *b* shows the maximum slop of 0.1714, followed *a* (0.0397),

and *c* has a minimum value (0.0156). The results are closely related to the matrix structure. As shown in Fig. 4b, the BiO<sub>8</sub> dodecahedra are ordered as layer-structure along *a*, *b* and *c* directions, connecting with RbO<sub>8</sub>, PO<sub>4</sub> and WO<sub>4</sub> to form the three-dimensional (3D) network. However, the inter-lamellar distances of BiO<sub>8</sub> layers are quite different along the different crystallographic orientation (Fig. 6). Along *c*-axis, the interlayer distance is as large as 10.23 Å (Fig. 6c), and it is only 3.53 Å (Fig. 6a) along *a*-axis and 2.09 Å (Fig. 6b) along *b*-axis which is the shortest value [32]. Therefore, with the minimum (Bi, Eu)O<sub>8</sub> layer distance, the cell parameter *b* is very sensitively to the variation of average ion size. This is another evidence as the Eu<sup>3+</sup> ions are successfully doped into the RBPWO host lattice.

#### 3.4. UV-visible DRS and photoluminescence properties

$$\alpha / S = F(R) = (1 - R)^2 / 2R \tag{2}$$

Where  $\alpha$ , *S* and *R* are the absorption, scattering and reflectance coefficients, respectively. As seen in **Fig. 7(a and c)**, the transitions in the DRS curve are the same as those of the PLE spectrum of RBPWO:0.9Eu<sup>3+</sup>. However, their relative intensities are very



**Fig. 5.** Cell parameters variation in the RBPWO: $xEu^{3+}$  (0.01  $\le x \le 1.0$ ) compounds.



**Fig. 6.** Crystal structure of RBPWO with the planar  $[BiO_8]_{\infty}$  layer along *a*, *b* and *c* directions.

different because the excitation transitions are also influenced by the QY of the correlative transitions [24]. Only the absorption generating efficient luminescence can be observed in the PLE spectrum. In the ultraviolet region (<350 nm), the very intense absorption bands are originated from the host lattice in DRS spectra. Whereas, the broad bands without fine structure are observed in PLE spectrum. As we all known, the Eu<sup>3+</sup> ion is inclined to receive one electron from neighbouring oxygen atom to achieve the half-filled electron configuration of Eu<sup>2+</sup>, and the according absorption transitions are called charge-transfer bands (CTB) [24,26,49]. Whereas, the CTB is inconspicuous in DRS curve for the overlap of the absorption of the host and the CTB of RBPWO:0.9Eu<sup>3+</sup>. In the NUV region (>350 nm), both DRS and PLE curves of the RBPWO:0.9Eu<sup>3+</sup> sample show several sharp absorption bands which can be assigned to the transitions from  ${}^{5}F_{0}$  to the excitation states of Eu<sup>3+</sup> ions. In the PLE spectrum, the  ${}^7F_0 \rightarrow {}^5L_6$ transition has the most intense absorption.

The room-temperature luminescence spectrum of RBPWO:Eu<sup>3+</sup> is the result of the radiative transitions from  ${}^{5}D_{0}$  level to the  ${}^{7}F$ manifold radiative transitions (**Fig. 7c**), and the  ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$  and  ${}^{5}D_{0} \rightarrow {}^{7}F_{0, 2, 4, 5}$  bands belong to the magnetic and forced electric dipole transitions, respectively [20,50–52]. As it is known, the  ${}^{5}D_{0}{}^{-7}F_{2}$  transition of Eu<sup>3+</sup> is responsible for the typical red luminescence emission, however, it is hypersensitive for the crystallographic surroundings [24,53]. In RBPWO:Eu<sup>3+</sup> compound, the low symmetry structure D<sub>2h</sub> of the host makes the  ${}^{5}D_{0}{}^{-7}F_{2}$  transition very intense and dominating the emission spectrum [32]. Meanwhile, the image (**Fig. 7a, inset**), the calculated emission CIE coordinates (0.6482, 0.3515) (**Fig. 7b**) and the high color purity (91%) further illustrate that the red-emitting material has been successfully obtained [54,55].

Additionally, as based on the Judd–Ofelt theory, the  ${}^{5}D_{0} \rightarrow {}^{7}F_{0}$ and  ${}^{5}D_{0} \rightarrow {}^{7}F_{3}$  transitions are forbidden. In the emission spectra,  ${}^{5}D_{0} \rightarrow {}^{7}F_{0}$  emission band is not observed, however, the weak



**Fig. 7.** UV-visible DRS of RBPWO and RBPWO:0.9Eu<sup>3+</sup> (inset shows the photograph of the RBPWO:0.9Eu<sup>3+</sup> tablet under the 365 nm excitation light) (a), the CIE color diagram (b) and the PLE and PL spectra of RBPWO:0.9Eu<sup>3+</sup> (c). The qualitative UV-visible DRS can be obtained according to the function (2) [48].

emission of  ${}^{5}D_{0} \rightarrow {}^{7}F_{3}$  is present in the range of 645–660 nm. The anomalous behavior can be explained by the *J*-mixing induced by the crystal-field perturbation [56]. The  ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$  transition is split into three sublevels due to the crystal field degeneracy [24]. The emission peaks at 589 nm and 594 nm are explicit, and, however, the one at 579 nm is very weak and almost invisible. Meanwhile, the  ${}^{5}D_{0} \rightarrow {}^{7}F_{4}$  emission is weaker than the  ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$  emission, but it is more intense than the  ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$  and  ${}^{5}D_{0} \rightarrow {}^{7}F_{3}$  transitions. Obviously, the  ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$  transition luminescence intensity

Obviously, the  ${}^{5}D_{0} \rightarrow {}^{7}F_{2}$  transition luminescence intensity (Fig. S2), the PL decay times (Fig. S3) and the AQY values (Fig. 8a, Table S3) all increase with the increasing of Eu<sup>3+</sup> content, and they reach the maximum values x = 0.9. Only when the Bi<sup>3+</sup> ions are

completely substituted by  $Eu^{3+}$  ions, the luminescence intensity decreased slightly. The critical distance ( $R_c$ ) between the adjacent  $Eu^{3+}$  ions can be calculated by the formula (3) [57]:

$$R_C \approx 2 \left(\frac{3V}{4\pi x_c N}\right)^{\frac{1}{3}}$$
(3)

Where *V*,  $x_c$  and *N* are the unit cell volume, critical quenching concentration and the number of Bi<sup>3+</sup> sites in the unit cell, respectively. In the structure of RBPWO, the unit cell volume is equal to 1810.69 Å and N = 8. When  $x_c$  = 0.9, the  $R_c$  value for concentration quenching is calculated to be 7.83 Å. It is very larger



**Fig. 8.** The absolute quantum yields of RBPWO: $xEu^{3+}$  (0 <  $x \le 1.0$ ) with 260 nm and 394 nm as excitation wavelength, respectively (a), and the coordination environment of  $Eu^{3+}$  (b).

than the minimal interlamellar distance between  $BiO_8$  layers (2.09 Å) (Fig. 6b). This result can be explained by the coordination environment of  $Eu^{3+}$ . As shown in Fig. 8b, each  $Eu^{3+}$  ion is surrounded by four PO<sub>4</sub> and two WO<sub>4</sub> tetrahedra. The adjacent luminescence activators are segregated by the above mentioned anions groups, which results in the very small interaction between europium ions and the concentration quenching is negligible [30,58,59]. Accordingly, the highest AQY values for the 394 and 260 nm excitation are 88.1% and 84.8%, respectively (Fig. 8a). The 394 nm is very close to the NUV output light of the commercial InGaN chip (395 nm).

#### 3.5. Thermal quenching properties

In general, the working current of modern w-LEDs is beyond 350 mA, and the internal temperature of the devices is about 423 K. To realizing the favorable color rendering index, the luminesce stability of the phosphors is essential. For the RBPWO:0.9Eu<sup>3+</sup> sample, the position of emission maximum peak keeps constant with increasing the environment temperature, which is beneficial for the stability of the color purity (Fig. 9a). Apparently, in higher temperature, the luminesce intensity of 614.5 nm red emission peak decreased gradually, and at 423 K, it decays to 79.7% of the initial value at 298 K (Table S4). As it was considered by Blasse, the charge-transfer state of  $Eu^{3+}$  had a significant impact on the luminescence thermal quenching process [33]. As depicted in Fig. 9b, the  ${}^{5}D_{0}$  ground and the CTB states of Eu ${}^{3+}$  are intersect at point S with the energy difference of E<sub>a</sub>. At room temperature, the transitions from the excited state  ${}^{5}D_{0}$  to the ground state  ${}^{7}F_{J}$ (I = 0-6) result in the emission spectrum. With the ambient temperature increases, when sufficient energy is provided by the thermal energy, the electrons located in  ${}^{5}D_{0}$  can jump to the CTB state. During this process, the electron must overcome the energy barrier E<sub>a</sub>, which is called activation energy. Then through the intersection S of CTB and  ${}^{7}F_{0}$ , the electrons located at the CTB state return to ground state directly. However, the transition is nonradiative and does not make contributions to the luminescence, which results in the quenching of luminescence. Based on the mechanism, it can be supposedly that, if the value of  $E_a$  is very high, the quenching of luminescence will occur at very high temperature. The activation energy  $E_a$  can be calculated from the following equation (4) [60].

$$I_{\rm T} = \frac{I_0}{1 + Ce^{-E_a/kT}} \tag{4}$$

Where  $I_0$  and  $I_T$  stand for the luminescence intensity at 614.5 nm at 298 K and a given temperature T, C is a constant, k is the Boltzmann constant (8.629 × 10<sup>-5</sup> eV K<sup>-1</sup>). As shown in Fig. S4, the ln[( $I_0/I_T$ ) –1] and 1/kT have a linear relationship, and the slope determines the activation energy  $E_a$  to be 0.23 eV.

Interestingly, while the emission intensity at 614.5 nm decrease, the emission lines were broaden due to the increasing of the electron-phonon interaction between the electrons and the lattice vibrations (Fig. 9a, inset) [61]. Due to the electron-phonon interaction, the emission lines blow 423 K in the RBPWO:0.9Eu<sup>3+</sup>, are broaden, which compensates the reduced emission intensity [61]. Therefore, the emission integrals are almost invariably within 423 K (within 1.5%) (Fig. S5 and Table S4), indicating that the total emission outputs change little with increasing temperature.

## 3.6. Electroluminescence of the fabricated w-LED device

To evaluate the potential application in solid lightings, a *w*-LED device was fabricated by coating the RBPWO:0.9Eu<sup>3+</sup> phosphor, the commercial blue BaMgAl<sub>10</sub>O<sub>17</sub>:Eu<sup>2+</sup> (BAM:Eu<sup>2+</sup>) and green (Sr,Ba)<sub>2</sub>SiO<sub>4</sub>:Eu<sup>2+</sup> phosphors on the surface of a NUV InGaN chip ( $\lambda$  = 395 nm). Under application of 20 mA (3.245 V) current, the white light with CIE color coordinates of (0.3842, 0.3908), T<sub>c</sub> of 4000 K and R<sub>a</sub> of 90 was obtained (Fig. 10a and b). The R<sub>a</sub> value is higher than that of the *w*-LED devices fabricated with Ca<sub>3</sub>Gd(AlO)<sub>3</sub>(BO<sub>3</sub>)<sub>4</sub>:Eu<sup>3+</sup> red phosphors (Ra = 72.5) [62]. Furthermore, the chromaticity parameters of the as-prepared *w*-LED operated in different currents were measured (Table S5), and little changes for the correlation parameters were observed with increasing the current.

## 4. Conclusions

The pure red-emitting RBPWO:Eu<sup>3+</sup> phosphors with high thermal stability and high quantum efficiency were synthesized. The RBPWO:Eu<sup>3+</sup> phosphors exhibit strong and broad CT absorption band centred at 260 nm and narrow absorption peak with



Fig. 9. (a) Variable-temperature photoluminescence spectra, and (b) Configuration coordinate diagram to show the process of thermal quenching of RBPWO:0.9Eu<sup>3+</sup>.



Fig. 10. (a) the electroluminescence spectrum (inset shows the fabricated device), and (b) the CIE chromaticity diagram with the coordinates of *w*-LED device fabricated with a 395 nm NUV chip.

maximum at 394 nm. The phosphors show intense red photoluminescence due to the  ${}^{5}\text{D}_{0}$ - ${}^{7}\text{F}_{2}$  transition dominating the emission spectra. Under the excitation at 394 nm, the AQY values can reach up to 88.1%. The total output lights almost invariable until the ambient temperature is higher than 423 K. The fabricated *w*-LED device can generate white light with CIE chromaticity coordinates of (0.3842, 0.3908), T<sub>c</sub> of 4000 K and R<sub>a</sub> of 90. Therefore, the RBPWO:Eu<sup>3+</sup> phosphors are expected as newly promising redemitting components for application in *w*-LEDs.

### **CRediT** authorship contribution statement

**Zhen Jia:** Data curation, Formal analysis, Writing - original draft. **Xiuling Zhang:** Methodology, Data curation. **Xiaoying Hua:** Methodology, Data curation. **Yan Dong:** Methodology, Data curation. **Hongliang Li:** Methodology, Data curation. **Chuanqi Feng:** Methodology, Data curation. **Yonggang Wang:** Formal analysis. **Mingjun Xia:** Methodology, Writing - review & editing.

## **Declaration of competing interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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#### Appendix A. Supplementary data

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