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An immutable array of TiO₂ nanotubes to pressures over 30 GPa

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Abstract

We report the successful formation of an immutable array of α -PbO₂ phase TiO₂ nanotubes by compression of a TiO₂ nanotube array in an anatase phase. During compression to 31.3 GPa, the TiO₂ nanotubes started to directly transform from an anatase phase to a baddeleyite phase at 14.5 GPa and completed the transition at 30.1 GPa. Under decompression, the baddeleyite phase transformed to an α -PbO₂ phase at 4.6 GPa, which was quenchable to ambient pressure. Notably the tubular array microstructure was retained after the application of ultra high pressure and undergoing a series of phase transformations. Measurements indicated that the nanotubes in the array possessed higher compressibility than in the bulk form. The highly aligned array structure is believed to reinforce the nanotubes themselves, giving exceptional stability. This, as well as the wall thickness, may also account for their different phase transition pathway.

Keywords: TiO₂, nanotube arrays, high pressure

(Some figures may appear in colour only in the online journal)

1. Introduction

Since the synthesis of carbon nanotubes in 1991 [1], efforts have focused on the preparation of solid material nano-scale tubular structures such as silica, boron nitride, gallium nitride, TiO₂, and V₂O₅ [2–8]. In the nanotube family, arrayed anodized TiO₂ nanotubes present significant advantages for applications in nano-devices, due to their oriented self-assembly on a substrate [9–14]. The anodized TiO₂ nanotube arrays are amorphous, which can transform to anatase and rutile phases through annealing at elevated temperatures [15]. However, the crystals in TiO₂ nanotubes grow quickly at high temperature and the tubular structure destructs above 800 °C. Therefore, a method that can transform TiO₂ nanotube arrays into other phases and maintain the assembled tubular morphology is highly desirable for various nano-device applications.

Pressure is an alternative method that may successfully tune materials' crystal structures and modify their physical and chemical properties. High-pressure studies on TiO₂ materials have so far revealed three crystalline phases in bulk TiO₂: the anatase phase, the α -PbO₂ phase (~2–5 GPa), and the baddeleyite phase (~12–15 GPa) [16–19]. The phase transformation pressures and stability of TiO₂ nanocrystals are size and morphology dependent. When the particle size is larger than 50 nm, they follow the same transition path as a bulk material. At sizes between 10–50 nm, they by pass the α -PbO₂ phase and directly transform to the baddeleyite phase [19, 20]. Compression of TiO₂ nanoparticles that are smaller than 10 nm leads to amorphization from their anatase phase [21–23]. The morphology-dependence of the anatase to baddeleyite phase transition at high pressure has been observed in 1D TiO₂ nanowires [24–26], 2D nanosheets [27] and random

nanotubes [28]. The phase transitions in either bulk or TiO₂ nanomaterials are irreversible. The high-pressure baddeleyite phase transforms to an α -PbO₂ phase after pressure release to ambient conditions. As a high-pressure phase of bulk TiO₂, α -PbO₂ is quenchable under decompression, before it transforms to a baddeleyite phase.

In the energy related applications such as solar cells [12], lithium batteries [29, 30], and photocatalysis [31, 32], TiO₂ is widely used due to its properties of charge separation and electron transport abilities. However, electron mobility in TiO₂ is very low ($0.1\text{--}4\text{ cm}^2\text{ V s}^{-1}$). Other than doping ion into TiO₂ materials to improve its properties, application of high pressure on TiO₂ materials is an effective method to reduce the resistivity [33]. Furthermore, the electrical conductivity of a quenched α -PbO₂ phase of TiO₂ has been enhanced by $\sim 40\%$ in comparison with that of the anatase phase. Thus, the α -PbO₂ phase of TiO₂ has the potential in the energy conversion applications.

TiO₂ nanotube arrays are promising candidates for various applications due to their unique, highly ordered arrangement and nanosized tubular structure. However, studies have only focused on their anatase and rutile phases. Further investigations of new crystalline structures besides these may provide new possibilities for TiO₂ nanotube array applications. Here, we present an investigation of phase transition behaviors in TiO₂ nanotube arrays under compression to over 30 GPa in a diamond anvil cell. Based on the characterization of the morphologies before and after compression, the TiO₂ nanotube arrays were ultra stable when subjected to high pressures over 30 GPa. We also discuss the size and morphology effects on phase transition behaviors and the stability of nanotube arrays at high pressures.

2. Experimental procedures

The sample was prepared using an electrochemical anodization method [13, 14]. A 0.25 mm thick, 99.5% pure titanium (Ti) foil (Alfa Aesar Co.) was used as our starting material and the substrate for the array synthesis. The foil was attached to the anode-working electrode. Another 0.025 mm thick, 99.9% purity platinum foil (Alfa Aesar Co.) was used as a counter-electrode. A solution of 98 vol% ethylene glycol (99.8% purity, JT Baker) with 2 vol% H₂O and an additional 0.25 wt% ammonium fluoride (NH₄F, 96% purity, Alfa Aesar Co.) served as an electrolyte during synthesis. A voltage of 40 V was applied for 2 h by a DC power supply (IT6874A, ITECH Co.) in the anodization process. The TiO₂ nanotube array film grew on the Ti film in the electrolyte solution during the process. The anodized sample was then washed with an ethanol solution and flushed with deionized water. Subsequently, it was annealed in an oven at 450 °C for 2 h to form the anatase phase of TiO₂ [14]. After annealing, the film was mechanically peeled off the Ti substrate for characterization and further experiments.

We compared the morphologies of the samples before and after compression with a JEM-2100F high-resolution transmission electron microscope (TEM) operated at 200 kV.

During high-pressure processing, a symmetric diamond anvil cell with a 400 μm diameter culet size was used to generate pressure. The sample was loaded into the sample chamber; a 120 μm diameter hole drilled in a preindented rhenium gasket. Subsequently, a ruby sphere as a pressure sensor and a mixture of methanol and ethanol with a volume ratio of 4:1 as a pressure transmitting medium (PTM) were also loaded into the sample chamber. *In situ* high-pressure Raman scattering and synchrotron x-ray diffraction (XRD) measurements were carried out during sample compression and decompression. The Raman measurement was performed by a Renishaw inVia Raman spectrometer with a laser wavelength of 532 nm. The *in situ* synchrotron XRD was undertaken at the 15U beam line of the Shanghai Synchrotron Radiation Facility. The x-ray wavelength was 0.6199 Å and the beam size was $10.6 \times 7.8\ \mu\text{m}^2$. The recorded 2D diffraction patterns were integrated into 1D profiles with the Fit2D program [34].

3. Results and discussion

Before high-pressure processing, the morphology and crystal structure of the sample were investigated by TEM and high-resolution TEM (HRTEM), as shown in figure 1. The sample had a highly ordered tubular array structure with a length of a few micrometers, tube diameter of ~ 100 nm and wall thickness of ~ 20 nm. The interplanar distance of 0.35 nm, which can be assigned to the (101) planes of anatase TiO₂, indicated that the nanotube array had a pristine anatase structure.

Figure 2(a) displays the XRD patterns in the compression process. At the lowest pressure of 1.3 GPa in the experiments, the diffraction peaks were indexed to the (101), (004) and (200) planes of the anatase phase. With increasing pressure, they slightly shifted to higher 2-theta angles, reflecting a reduction of the interplanar distance with compression. Two new peaks at 12.4° and 13.7° appeared at 16.5 GPa, along with the disappearance of the (004) peak of the anatase phase. These peaks were assigned to the ($\bar{1}11$) and (111) peaks of the baddeleyite phase. All peaks of the anatase phase disappeared and the peaks of baddeleyite phase remained at the experiment's highest pressure of 30.1 GPa. Thus, based on these measurements, the transition from the anatase to baddeleyite phase started from 16.5 GPa and completed at 30.1 GPa. Figure 2(b) displays the diffraction spectra during decompression. A regular drift of all peaks to lower angles was observed during pressure reduction until 9.2 GPa. However, when the sample was finally quenched to ambient conditions from 9.2 GPa, only two diffraction peaks could be indexed to the (111) and (110) peaks of the α -PbO₂ phase. Thus, the sample eventually transformed to and maintained the α -PbO₂ phase at ambient pressure.

To verify the observed phase transition sequence of the TiO₂ nanotube arrays from XRD measurements, Raman spectra were applied to characterize the pressure-induced structural phase transition. At ambient conditions, the anatase TiO₂ tetragonal lattice belonged to the space group D_{4h}^{19} (I4/

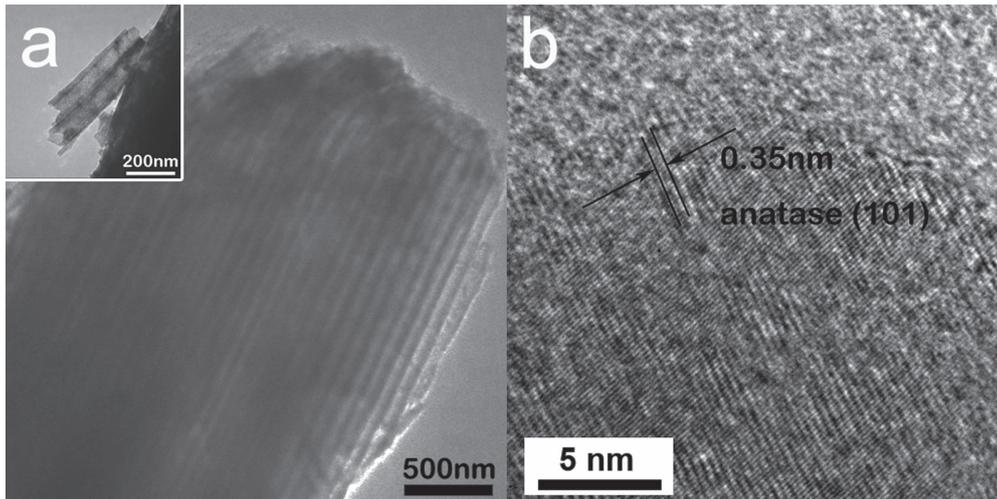


Figure 1. TEM (a) and HRTEM (b) images of the initial sample.

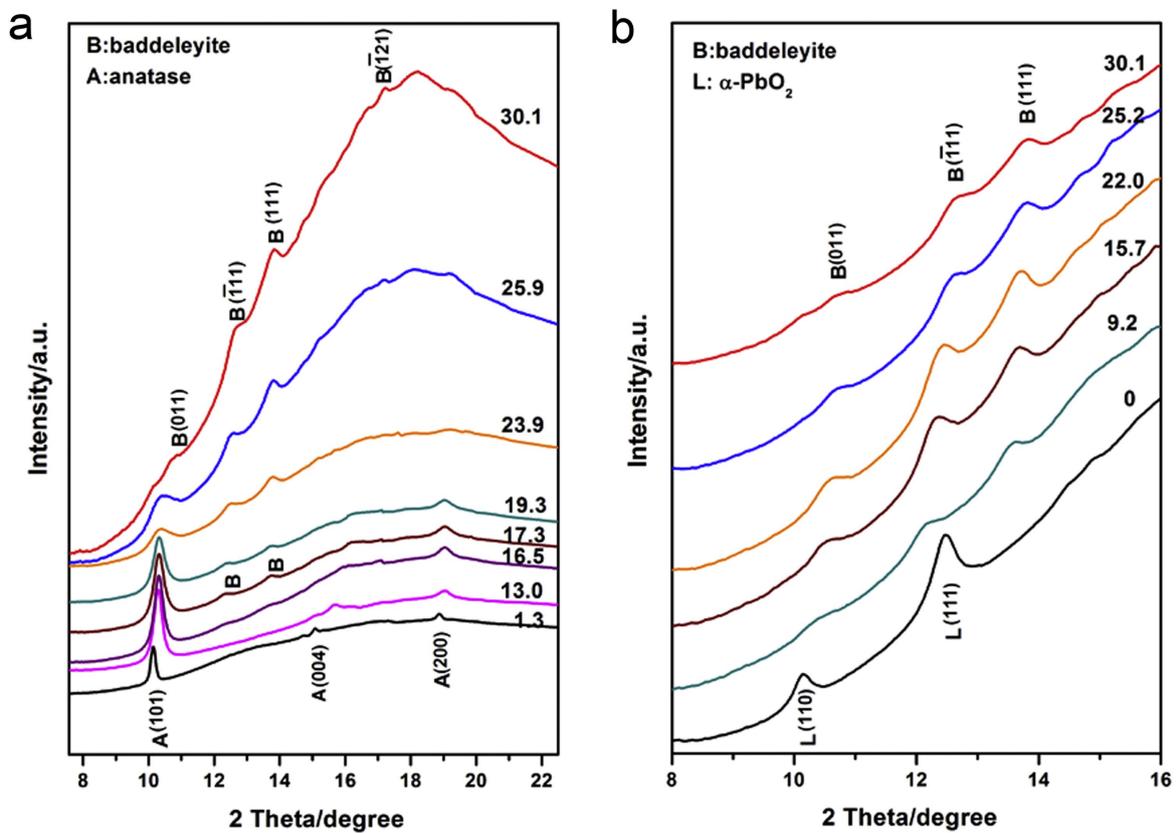


Figure 2. X-ray diffraction patterns of TiO_2 nanotube arrays at selected pressures during: (a) compression and (b) decompression. The pressure in GPa is labeled above each curve.

amd) [35]. According to its primitive unit cell and factor group analysis, it consisted of six Raman modes including the $3E_g$, $2B_{1g}$ and A_{1g} modes. Figure 3 shows the Raman spectra at selected pressures. The spectrum at ambient pressure (the bottom spectrum in figure 3(a)) shows Raman modes at 142, 194, 394, 515 and 635 cm^{-1} , which can be identified as $E_{g(1)}$, $E_{g(2)}$, $B_{1g(1)}$, $A_{1g(1)} + B_{1g(2)}$, and $E_{g(3)}$ [36]. Upon compression, all the Raman modes showed routine blue-shift along with a significant intensity decrease and

width increase. A new peak at 489 cm^{-1} was identified from the baddeleyite structure [22, 37] at 14.5 GPa, indicating that a phase transition from the anatase phase to a baddeleyite phase had occurred, which is consistent with our XRD observation. At pressures above 14.5 GPa, this peak grew and another peak from the baddeleyite phase was observed at $\sim 230 \text{ cm}^{-1}$ when pressure reached 18.6 GPa. At 31.3 GPa, all the peaks from the anatase phase disappeared and only baddeleyite peaks were visible. It indicated that the anatase phase

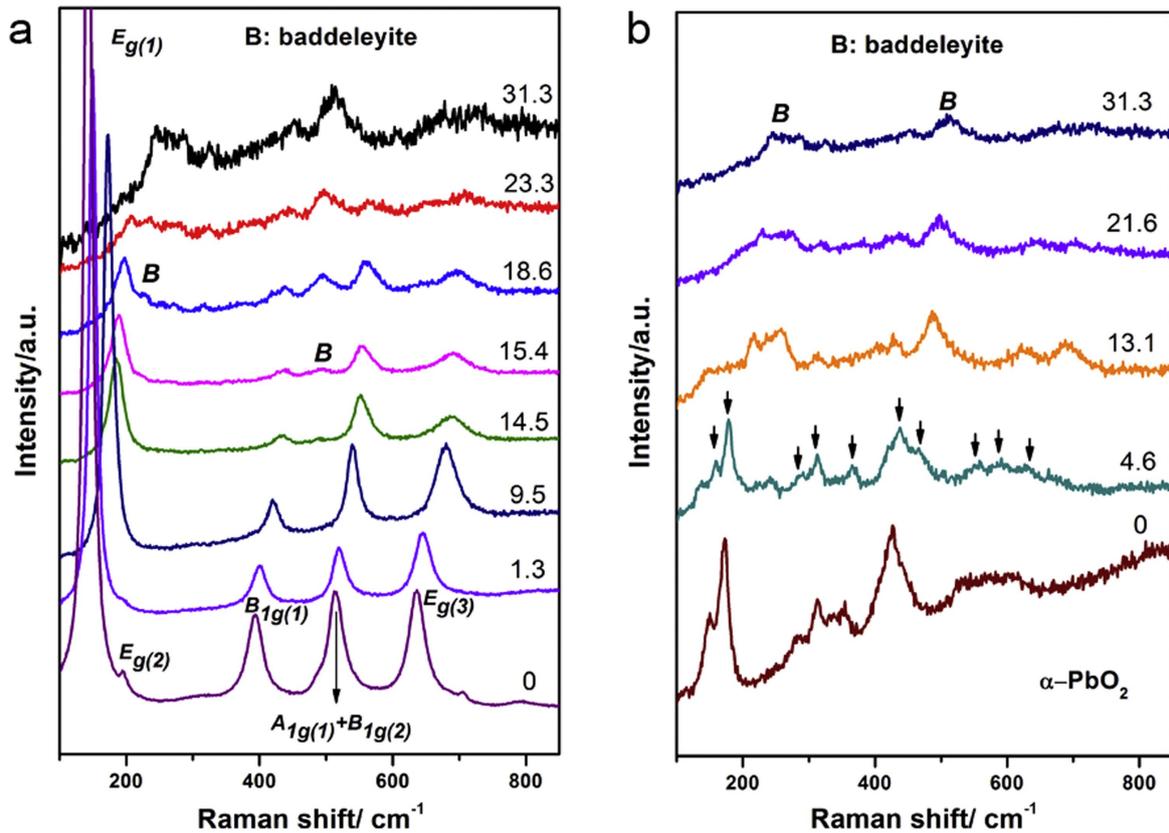


Figure 3. Raman spectra of TiO₂ nanotube arrays during: (a) compression and (b) decompression. The pressure in GPa is labeled above each spectrum.

had completely transformed to the baddeleyite phase. During decompression, all the peaks from the baddeleyite phase remained until 13.1 GPa, as shown in figure 3(b). When pressure decreased to 4.6 GPa, new peaks were observed at 160, 179, 287, 313, 367, 436, 554, 588, and 632 cm⁻¹. These peaks were all ascribed to scattering from the α -PbO₂ phase. When the sample was finally quenched down to ambient pressure, all the Raman modes from the baddeleyite phase disappeared. The α -PbO₂ phase remained and the intensity of those Raman peaks significantly increased. Thus, we determined that the baddeleyite phase transformed to the α -PbO₂ phase during decompression at 4.6 GPa and the α -PbO₂ phase was quenchable to ambient pressure, which agrees with our XRD results.

To investigate the morphological stability of the TiO₂ nanotube arrays, as well as the crystal structure change, TEM and HRTEM measurements were performed on the quenched sample, as shown in figure 4. The tubular array structure survived after undergoing both compression and decompression (figure 4(a)). Figure 4(b) shows good crystallinity of the quenched TiO₂ nanotube and a clear interplanar distance of 0.28 nm, which is ascribed to the (111) planes of the α -PbO₂ phase. This is consistent with the x-ray and Raman measurement results and demonstrates that the TiO₂ nanotube arrays undergo an anatase to baddeleyite then to α -PbO₂ phase transitions during compression up to ~31 GPa and decompression down to ambient conditions, respectively, without any morphology damage.

The XRD and Raman results clearly suggest that the high-pressure behavior of the anatase structure TiO₂ nanotube arrays differ from the bulk material. It is well known that bulk TiO₂ has a phase transition path from an anatase to an α -PbO₂ phase, and then to a baddeleyite phase during compression. The TiO₂ nanotube arrays present a phase transition sequence similar to some of the nano-sized TiO₂ materials listed in table 1. For the nanoparticles, the phase transition paths are size-dependent. There are two critical sizes (10 and 50 nm) for the anatase TiO₂ particles' transformation to different high-pressure phases. The other TiO₂ materials (nanowires, nanosheets, and random nanotubes) present morphology-dependent phase transitions. Our TiO₂ material has tubular morphology, which allows the PTM to penetrate the tubes and provide a hydrostatic environment for the tube walls. The wall thickness of the tubes is about 20 nm, which is a favorable dimension for the direct phase transition from the anatase phase to the baddeleyite phase. Therefore, we suggest that the unique tubular structure, as well as the nanosized wall thickness, play determining roles in the phase transition behaviors of the TiO₂ nanotube arrays. As shown by our TEM studies, the ultra-stable tubular array structure of this TiO₂ material may also result from its special morphology with a perfectly aligned and compact arrangement. The self-assembled array of TiO₂ nanotubes grew due to the competitive reactions between anodic oxidation and chemical dissolution in the electrolyte [38]. By controlling the fluoride concentration, temperature, anodization potential and reaction

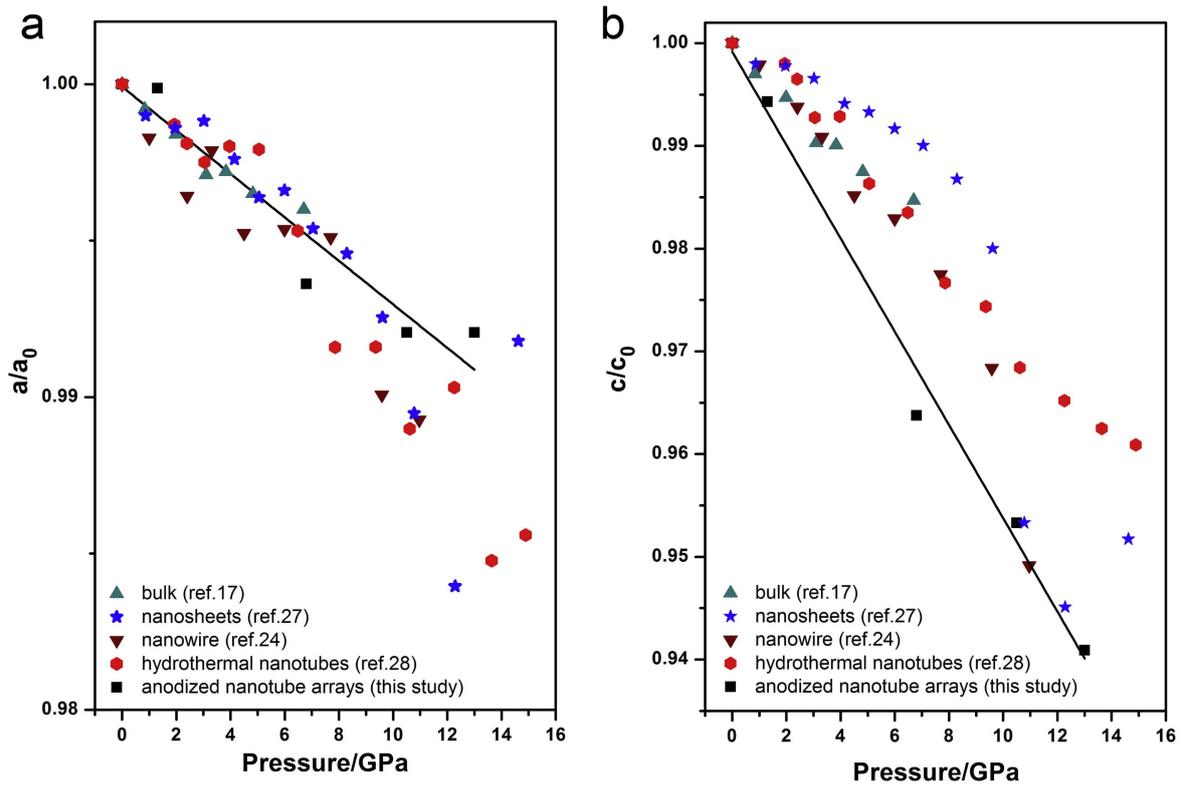


Figure 6. Normalized cell parameters of TiO₂ nanotube arrays at pressures. Solid lines are the linear fitting in this study.

Analysis of the volume reduction as a function of pressure allows further understanding of TiO₂ nanotube array compressibility. This is compared with earlier studies of TiO₂ with various sizes and morphologies in figure 7. The bulk modulus is calculated by fitting the third-order Birch Murnaghan equation:

$$P = 3/2B_0[(V/V_0)^{-7/3} - (V/V_0)^{-5/3}] \times \{1 + 3/4(B'_0 - 4)[(V/V_0)^{-2/3} - 1]\},$$

where V is the volume at pressure P , V_0 is the volume at zero-pressure, and B_0 and B'_0 are the isothermal bulk modulus and its pressure derivative, respectively. The bulk modulus (B_0) of the nanotube arrays is determined as 140 ± 9 GPa when B'_0 is fixed at 4. It is smaller than those (166–317 GPa) of the reported bulk TiO₂ [17] and nano-sized materials [20, 21, 24, 27, 28, 40, 41]. Compared to the compressibility of different morphologies along the a - and c -axes, the slope of c/c_0 from the nanotube arrays was two times that of the bulk. We ascribe this low bulk modulus in the nanotube arrays to the high shrinking rate in the c -axis, which is shown in figure 6.

4. Conclusions

We studied the behaviors of anatase TiO₂ nanotube arrays at high pressures. Upon compression, the intermediate α -PbO₂ phase was missing, which contrasts with the phase transitions in a bulk TiO₂ material. The anatase phase started transformation to a baddeleyite phase at

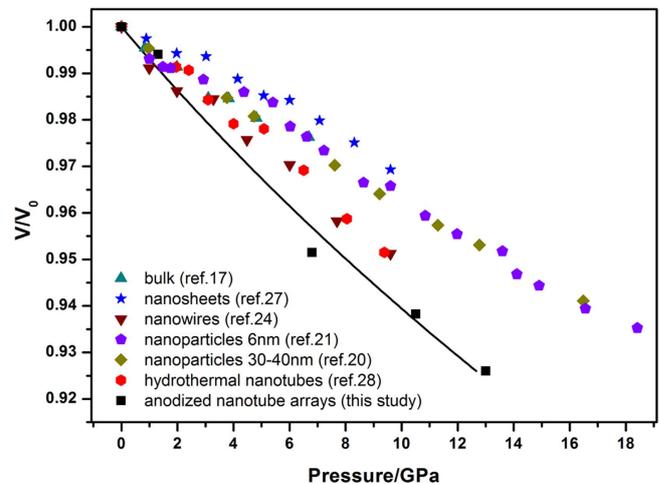


Figure 7. The volume reduction of the anatase phase of TiO₂ nanotube arrays.

14.5 GPa upon compression. The baddeleyite phase transformed to the α -PbO₂ phase during decompression. The 3D TiO₂ nanotubes had ultra-stable morphology and presented very high compressibility under pressure up to 31.3 GPa, which demonstrates that TiO₂ nanotubes with a good crystalline α -PbO₂ phase can be formed by high-pressure processing. These pressure-induced phase transitions in TiO₂ nanotube arrays provide a new candidate to develop α -PbO₂ phase based nano-devices of nanotubes.

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